

## A Refinement of the Crystal Structure of Gypsum $\text{CaSO}_4 \cdot 2\text{H}_2\text{O}$

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With the use of three-dimensional X-ray diffraction data the crystal structure of gypsum has been refined by least-squares methods. Initial parameters used were those found by Atoji & Rundle [*J. Chem. Phys.* (1958), **29**, 1306] from neutron-diffraction data. Significant changes were made in the  $x$  and  $z$  parameters of the oxygen atoms of the sulphate ion and the oxygen atom of the water molecule. Bond lengths around the sulphate ion show equal sulphur-to-oxygen distances of mean value 1.459 Å, but one pair of oxygen-to-oxygen distances of 2.323 Å and the corresponding pair of oxygen-sulphur-oxygen bond angles of 105.5° are significantly different from the others, which have mean values of 2.411 Å and 111.5°, respectively. As a result the symmetry of the sulphate ion is lowered from a tetrahedron to a sphenoid. The new oxygen parameters indicate that there are two distinctly different hydrogen-bond lengths of 2.816 and 2.896 Å rather than one only of mean value 2.820 Å as found by Atoji & Rundle. A re-working of the data of Atoji & Rundle produced refinement beyond that which they reported and clearly indicated the two different hydrogen-bond lengths (difference  $0.065 \pm 0.018$  Å) shown by the X-ray work. The re-refining of the neutron data also confirmed the lowering of the symmetry of the sulphate ion, although this did not show up in Atoji & Rundle's refinement. On the basis of the new parameters for the hydrogen atoms from the re-worked neutron data and the parameters for the water oxygen found from the X-ray data, the two O-H distances in the water molecule are calculated as 0.962 and 0.944 Å but the difference is not statistically significant. The details of the refined structure are in agreement with conclusions drawn by other workers from spectroscopic data. The anisotropic thermal parameters are discussed with respect to the structure.

### Introduction

The crystal structure of gypsum,  $\text{CaSO}_4 \cdot 2\text{H}_2\text{O}$  was determined by Wooster (1936) from two-dimensional X-ray diffraction data with Fourier methods, which gave an  $R$  index of 13.6% on observed data. It was further refined by Atoji & Rundle (1958) from  $0kl$ ,  $hk0$  and  $hkh$  neutron-diffraction data by means of both Fourier and least-squares refinement methods to an  $R$  index of 15.5% on observed data. The neutron-diffraction work permitted the fixing of the positions of the hydrogen atoms and indicated small but significant changes in the positions of all of the oxygen atoms from those established by Wooster. Denne & Jones (1969) re-examined the neutron-diffraction data of Atoji & Rundle by an occupation-factor approach and found that the distribution of the hydrogen atoms was centrosymmetric.

By modern standards the crystal structure of gypsum is only poorly refined and the present study was undertaken to refine the structure of this important industrial and building material by using three-dimensional X-ray data collected on a modern semi-automatic counter diffractometer. The study has shown (Cole & Lancucki, 1972*a, b*) that the symmetry of the  $\text{SO}_4$  ion in gypsum is lower than tetrahedral, confirming the earlier deductions of Rousset & Lochet (1945) from Raman, and Hass & Sutherland (1956) from infrared spectra. It has also been found (Cole & Lancucki, 1973) that

the water molecule is asymmetrical, confirming the deductions of Seidl, Knop & Falk (1969) from infrared spectra of partially deuterated gypsum.

### Experimental

The gypsum used was from Monte San Derato Bologna, Italy. It had the following composition: CaO 32.5,  $\text{SO}_3$  46.4,  $\text{H}_2\text{O}$  21.0,  $\text{SiO}_2$  + other insolubles 0.01,  $\text{R}_2\text{O}_3 < 0.01$ ,  $\text{MgO} < 0.005$ ,  $\text{Na}_2\text{O} < 0.01$ ,  $\text{SrO}$  0.025,  $\text{BaO} < 0.005$ ,  $\text{CO}_2 < 0.01$ , and Cl 0.01%. It was shaped into a sphere of about 0.3 mm diameter, and ellipsoids of 0.3 mm on the minor axis and 0.37 mm on the major axis as described by Cole & Lancucki (1965). This involved spinning cut cubes of the crystals to ellipsoids by circulating kerosine in a Bond cell, and hand shaping the ellipsoids to spheres, followed by etching in running water. The crystals used in the analysis showed some surface streaking for many low-angle reflexions.

The unit-cell dimensions were determined by the  $\theta$  method, as described by Weiss, Cochran & Cole (1948), with a Unicam single-crystal goniometer. The values of  $a^*$ ,  $b^*$ ,  $c^*$  and  $\beta^*$  were obtained directly from  $h00$ ,  $0k0$  and  $00l$  reflexions by extrapolating the regression equations on  $\cos^2 \theta$  to  $\theta = 90^\circ$ . The new unit-cell parameters and those found by Atoji & Rundle are reported below together with other relevant crystallographic data relating to the Bragg cell (Bragg, 1937).

	This study	Atoji & Rundle
Space group	<i>I</i> 2/ <i>a</i>	<i>I</i> 2/ <i>a</i>
<i>a</i>	5·670 ± 0·002 Å	5·68 ± 0·008 Å
<i>b</i>	15·201 ± 0·002	15·18 ± 0·009
<i>c</i>	6·533 ± 0·002	6·52 ± 0·008
$\beta$	118° 36' ± 4'	118° 23' ± 20'
Cell volume	494·37 Å <sup>3</sup>	494·59 Å <sup>3</sup>
Cell contents	4(CaSO <sub>4</sub> · 2H <sub>2</sub> O)	4(CaSO <sub>4</sub> · 2H <sub>2</sub> O)

The spherically shaped crystal was rotated about the *b* axis and the ellipsoidal crystal was rotated about the *a* axis and initially Weissenberg photographs were taken with Mo *K* $\alpha$  radiation and multiple-film packs with Sn foil (0·025 mm thick) as an absorber. The intensities of the reflexions were estimated by eye with a multiple-film scale prepared under the same conditions as for the films. Of these data use was made of only the *b*-axis intensities and these were corrected for both Lorentz-polarization and absorption. The latter correction was small since  $\mu r = 0\cdot54$ . Data from the same crystals rotated about the same axes were subsequently collected on a computer-controlled Buerger-Supper equi-inclination X-ray diffractometer as described by Freeman, Guss, Nockolds, Page & Webster (1970). A constant slow scan speed of 0·05° sec<sup>-1</sup> was used with filtered Mo *K* $\alpha$  radiation, although some data about the *b* axis were collected with radiation monochromatized with a plane graphite crystal. For each level of reciprocal space the varying indices were allowed to take both positive and negative values so that reflexions with a multiplicity of two were measured twice and those with a multiplicity of four were measured four times.

The counter-diffractometer data were processed in the normal manner for application of Lorentz-polarization factors and calculations of  $I_{hkl}$ ,  $\sigma(I_{hkl})$ ,  $F_{hkl}$  and  $\sigma(F_{hkl})$ . Because of streaking, many low-angle reflexions suffered a background imbalance, *i.e.* they did not satisfy the test  $|B_1 - B_2| < (BI) \cdot \sigma(B) + (IC)$  where  $\sigma(B) = \sqrt{(B_1 + B_2)}$ ,  $BI = 10\cdot0$  and  $IC = 0\cdot01$ . In these instances the lowest background value was assumed for both  $B_1$  and  $B_2$  and the data processed in the normal way. In the merging of the two sets of data from *a* and *b* axes all levels were scaled manually to the  $2kl$  level on the basis of the sums of the *F* values of common reflexions, thereby yielding finally only one scale factor (*Sq*) between observed and calculated *F* values. This was necessitated by the large amount of data being handled compared with the memory limitations of available computer programs. 2193 non-equivalent reflexions were measured as present and a further 2158 were recorded as being too weak to be observed ( $I < I_{min} = RIM \cdot \sigma(B)/\sqrt{2}$  where  $RIM = 3\cdot5$ ). For these *I* was set equal to 0·33 *I*<sub>min</sub> and  $\sigma(I)$  was set equal to 0·29 *I*<sub>min</sub>.

**Structure analysis**

A least-squares refinement was made on the counter-diffractometer data using the Busing, Martin & Levy

Table 1. Isotropic temperature coefficients (*B*), fractional atomic parameters (*x/a*, *y/b*, *z/c*) from X-ray film and counter data of the present work and the neutron counter data of Atoji & Rundle (1958), as refined by them, and as re-refined in the present work

	<i>B</i> (Å <sup>2</sup> )				<i>x/a</i>				<i>y/b</i>				<i>z/c</i>			
	X-ray		Neutron		X-ray		Neutron		X-ray		Neutron		X-ray		Neutron	
	Film	Counter	Counter	Re-refined	Film	Counter	Counter	Re-refined	Film	Counter	Counter	Re-refined	Film	Counter	Counter	Re-refined
Ca	0·88 (5)	0·51 (1)	1·66	1·72 (25)	0·5000 (0)	0·5000 (0)	0·5000 (0)	0·5000 (0)	0·5000 (0)	0·5000 (0)	0·5000 (0)	0·5000 (0)	0·2500 (0)	0·2500 (0)	0·2500 (0)	0·2500 (0)
S	0·66 (6)	0·63 (1)	1·84	1·40 (36)	0·0000 (0)	0·0000 (0)	0·0000 (0)	0·0000 (0)	0·7500 (0)	0·7500 (0)	0·7500 (0)	0·7500 (0)	0·7500 (0)	0·7500 (0)	0·7500 (0)	0·7500 (0)
O(1)	1·49 (7)	1·24 (2)	2·13	1·67 (17)	0·9623 (10)	0·9607 (3)	0·9665 (11)	0·9665 (11)	0·5515 (9)	0·5518 (2)	0·5518 (2)	0·5518 (2)	0·5515 (9)	0·5518 (2)	0·5495 (16)	0·5495 (16)
O(11)	1·41 (6)	1·18 (2)	2·18	1·94 (18)	0·7586 (10)	0·7592 (2)	0·7586 (10)	0·7592 (2)	0·6674 (8)	0·6661 (2)	0·6661 (2)	0·6674 (8)	0·6735 (14)	0·6674 (16)	0·6674 (16)	0·6674 (16)
O(W)	2·19 (10)	1·90 (3)	2·41	2·50 (20)	0·3784 (12)	0·3778 (3)	0·3784 (12)	0·3778 (3)	0·1826 (6)	0·1826 (6)	0·1826 (6)	0·1826 (6)	0·4597 (11)	0·4588 (2)	0·4553 (20)	0·4553 (20)
H(1)			3·82	4·05 (33)			4·05 (33)	4·05 (33)				0·1601 (10)			0·4974 (50)	0·4974 (50)
H(2)			3·56	3·98 (35)			3·98 (35)	3·98 (35)				0·2430 (12)			0·4880 (40)	0·4880 (40)

(Numbers in parentheses are standard deviations in the least significant digits.)

(1962) computer program *ORFLS* with isotropic temperature factors and the initial atomic parameters and temperature factors found by Atoji & Rundle (1958). The  $f$  curves for  $S^{6+}$ , O, and H were taken from Mirkin (1964),  $Ca^{2+}$  from *International Tables for X-ray Crystallography* (1962) and  $O^{2-}$  from the example on quartz featured in *ORFLS*. With structure factors having unit weights ( $w$ ) and  $hkl$  reflexions for  $k$  between 0 and 3, an overall disagreement index  $R = \sqrt{\sum w(F_o - SqF_c)^2} / \sqrt{\sum wF_o^2} = 0.089$  was achieved, excluding hydrogen atoms, after four cycles of refinement for 365 non-equivalent reflexions. With anisotropic temperature factors,  $R = 0.074$ . For the full data set of 2193 non-equivalent reflexions refined with the use of a weighting function  $w = [\sigma(F_o)]^{-2}$  an  $R$  value = 0.055 was obtained with isotropic temperature factors after three cycles of refinement; with anisotropic temperature factors,  $R = 0.048$  after three cycles of refinement. When hydrogen atoms were introduced into the isotropic refinements with the parameters found by Atoji & Rundle,  $R$  became 0.059. Attempts to refine the hydrogen parameters with the full data set were unsuccessful, the temperature factors assuming unrealistic values. Some refinement of the positional parameters for hydrogen was achieved by using reflexions with  $\sin \theta < 0.4$  and keeping the thermal parameters constant, as recommended by Huber-Buser (1971). The final parameters gave values of bond lengths within the water molecule considerably different from acceptable values. Difference Fourier maps based on  $\Delta F = (F_o - SqF_c)$  and excluding hydrogen atoms from the calculated structure factors were computed in the form of projected sections in the vicinity of the hydrogen atoms for both sets of data. Weak electron density maxima were found in the vicinity of the expected values and the positional parameters deduced from the peaks [for H(1)  $x/a = 0.2500$ ,  $y/b = 0.1463$ ,  $z/c = 0.4680$ ; for H(2)  $x/a = 0.3917$ ,  $y/b = 0.2166$ ,  $z/c = 0.5033$ ] were in closer agreement with the least-squares refinement than with the values found by Atoji & Rundle from neutron-diffraction data. The hydrogen positions are discussed further in a later paragraph.

The positional and thermal parameters for O(1), O(11), O( $W$ ), Ca and S are given in Table 1 with their standard deviations for the isotropic refinement of the 2193 non-equivalent reflexions. They are compared

Table 2. *Refined anisotropic thermal parameters* ( $\times 10^4$ )

The anisotropic temperature factor has the form

$$T = \exp [-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2kl\beta_{23} + 2hl\beta_{13})].$$

(Numbers in parentheses are standard deviations in the least significant digit.)

	$\beta_{11}$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$
Ca	69 (1)	9 (0)	59 (1)	0	23 (1)	0
S	52 (2)	7 (0)	46 (1)	0	13 (1)	0
O(1)	115 (4)	15 (0)	89 (3)	-3 (1)	42 (3)	4 (1)
O(11)	81 (4)	14 (0)	92 (3)	-8 (1)	15 (3)	4 (1)
O( $W$ )	225 (6)	16 (1)	204 (5)	3 (2)	149 (5)	-7 (1)

with the previous values found by Atoji & Rundle (1958). The thermal parameters arising from the anisotropic refinement of the same 2193 reflexions are given in Table 2.

Because the positional parameters for O(1), O(11) and O( $W$ ) differed significantly from those found by Atoji & Rundle using neutron methods, the film data collected from the rotation about the  $b$  axis were also subjected to a least-squares refinement. In this case each level of reciprocal space was scaled independently to the calculated structure factors and the scale factors and atomic parameters were refined with isotropic temperature factors to an  $R$  value of 0.011 with a unit weighting function. The parameters (Table 1) showed similar differences from the neutron-diffraction data reported by Atoji & Rundle as occurred with the X-ray counter diffraction data.

In view of these differences a re-refinement was made of the neutron-diffraction data with the computer program *ORFLS* using isotropic temperature factors and initial atomic parameters and temperature factors as found by Atoji & Rundle. The scattering lengths of 0.49, 0.31, 0.58 and -0.38 in Fermi units ( $10^{-11}$  mm) for Ca, S, O, and H, respectively were the same as those used by Atoji & Rundle. In four cycles of refinement  $R$  was lowered from 0.154 to 0.092, to give atomic parameters and temperature factors nearly identical with those in Table 1. The values in Table 1 relate to refinements made with the more recent scattering lengths of 0.490, 0.280, 0.577, and -0.378 for Ca, S, O, and H, respectively. These gave a minimum  $R$  value of 0.091. The new scattering lengths mainly influenced the temperature factors, particularly that for S. The atomic parameters for O(1), O(11) and O( $W$ ) have approached those found from the X-ray work and H(1) and H(2) are in positions considerably different from those found by Atoji & Rundle. Attempts to refine the neutron-diffraction data anisotropically were only partially successful. Atoms O(1), Ca and S were not positive-definite after the first cycle and if these were left isotropic, atom H(1) became not positive-definite after the second cycle of refinement. The  $R$  index was lowered to a minimum value of 0.081.

The observed and calculated X-ray structure factors based on the anisotropic X-ray refinement parameters for Ca, S, O(1), O(11) and O( $W$ ) and the isotropic neutron re-refined parameters for hydrogen are given in Table 3. The observed and calculated neutron structure factors based on the re-refined isotropic parameters of Atoji & Rundle's observed data are given in Table 4.

## Discussion

In the crystal structure of gypsum (Wooster, 1936; Atoji & Rundle, 1958) there are two crystallographically distinct sulphate oxygen atoms [O(1) and O(11)]. The new refinements have produced significant changes in the  $x$  and  $z$  parameters (Table 1) of the oxygen atoms O(1) and O(11) and the oxygen atom O( $W$ ) of the water



Table 3 (cont.)

A large table of numerical data with multiple columns and rows, likely representing crystallographic parameters or diffraction data.

ion having a mean value of 2.411 Å, and the corresponding angles a mean value of 111.5°. The sulphur-oxygen distances are essentially identical at 1.459 Å and can be compared with an average distance of 1.49 Å in other inorganic structures (International Tables for X-ray Crystallography, 1962). The new values of bond lengths and angles for the sulphate ion

in gypsum lower its symmetry from a tetrahedron in the point group 43m (hexatetrahedral) to a sphenoid in the point group 42m (tetragonal bisphenoidal). The fact that the symmetry of the sulphate ion must be lower than tetrahedral has long been established from both Raman (Roussset & Lochet, 1945) and infrared spectra (Hass & Sutherland, 1956), both of which show

Table 4. *Observed and calculated neutron structure factors for  $\text{CaSO}_4 \cdot 2\text{H}_2\text{O}$* 

Based on an isotropic re-refinement of Atoji & Rundle's observed data. Structure factors are scaled to  $100F$  in Fermi units, where  $F$  is calculated for one  $\text{CaSO}_4 \cdot 2\text{H}_2\text{O}$ .

$h$	$k$	$l$	FO	SQFC	$h$	$k$	$l$	FO	SQFC
2	0	0	52.9	53.2	0	17	1	19.1	19.2
4	0	0	61.4	52.8	0	2	2	28.1	-31.1
6	0	0	9.7	12.5	0	4	2	61.6	-61.8
0	2	0	81.7	78.4	0	6	2	59.3	56.9
0	4	0	44.4	-37.5	0	8	2	2.3	-6.0
0	6	0	61.4	54.2	0	10	2	27.5	23.1
0	8	0	21.0	-20.9	0	12	2	57.0	-66.8
0	10	0	51.3	49.7	0	14	2	10.2	-8.7
0	12	0	32.4	-31.5	0	16	2	47.9	46.9
0	14	0	15.6	13.6	0	1	3	5.0	7.7
0	16	0	31.9	30.0	0	3	3	34.7	34.9
0	18	0	66.2	-65.7	0	5	3	5.9	-7.0
0	0	2	31.7	-37.6	0	7	3	51.3	-49.7
0	0	4	8.6	-2.2	0	9	3	49.6	48.1
0	0	6	25.5	-26.6	0	11	3	3.8	-2.1
1	1	0	16.3	14.3	0	13	3	37.7	-39.6
1	3	0	6.4	-5.4	0	15	3	22.9	19.5
1	5	0	44.3	-40.2	0	2	4	11.3	13.4
1	7	0	52.3	50.0	0	4	4	49.6	-49.0
1	9	0	47.4	44.0	0	6	4	26.3	-23.4
1	11	0	81.3	-78.0	0	8	4	39.9	-40.1
1	13	0	15.2	10.3	0	10	4	26.7	25.3
1	15	0	24.9	26.2	0	12	4	< 2.0	-2.0
1	17	0	6.5	-8.1	0	1	5	< 2.0	-2.1
2	2	0	39.4	-39.0	0	3	5	11.8	8.5
2	4	0	125.1	-118.2	0	5	5	36.5	-34.8
2	6	0	16.9	-17.5	0	7	5	70.0	-67.7
2	8	0	21.0	-13.6	0	9	5	33.5	35.3
2	10	0	23.7	-29.8	0	11	5	15.8	-17.8
2	12	0	13.4	9.5	0	2	6	35.1	35.3
2	14	0	53.1	52.4	0	4	6	43.3	40.4
2	16	0	27.1	27.2	2	0	2	48.2	40.1
3	1	0	24.3	20.9	2	0	4	45.1	27.9
3	3	0	64.2	-63.3	1	2	$\bar{1}$	70.2	-75.9
3	5	0	3.0	3.4	2	2	$\bar{2}$	18.4	-21.9
3	7	0	29.1	31.2	3	2	$\bar{3}$	23.8	19.4
3	9	0	10.6	-8.7	4	2	$\bar{4}$	< 4.0	-5.9
3	11	0	11.3	-8.8	5	2	$\bar{5}$	42.1	-35.7
3	13	0	23.3	-21.9	1	4	$\bar{1}$	33.2	-36.5
3	15	0	20.6	24.9	2	4	$\bar{2}$	25.8	-24.9
4	2	0	70.2	62.8	3	4	$\bar{3}$	43.6	-40.4
4	4	0	25.6	26.4	4	4	$\bar{4}$	21.0	14.4
4	6	0	50.3	-49.4	5	4	$\bar{5}$	34.2	-31.0
4	8	0	63.6	60.8	1	6	$\bar{1}$	15.4	13.0
4	10	0	11.6	-12.3	2	6	$\bar{2}$	90.0	102.7
4	12	0	9.5	-8.6	3	6	$\bar{3}$	28.2	27.7
4	14	0	32.3	31.5	4	6	$\bar{4}$	86.7	-86.0
5	1	0	22.0	22.4	5	6	$\bar{5}$	15.9	16.9
5	3	0	21.8	-19.6	1	8	$\bar{1}$	23.3	-24.1
5	5	0	21.4	18.3	3	8	$\bar{3}$	31.1	-36.2
5	7	0	21.6	20.6	4	8	$\bar{4}$	23.8	-22.8
5	9	0	< 4.0	-4.6	1	10	$\bar{1}$	23.9	-22.6
5	11	0	8.6	4.8	2	10	$\bar{2}$	14.3	-11.6
6	2	0	21.6	-19.8	3	10	$\bar{3}$	120.1	-139.3
6	4	0	32.6	-32.8	4	10	$\bar{4}$	11.3	6.0
0	1	1	< 5.0	3.1	1	12	$\bar{1}$	26.4	29.4
0	3	1	20.8	20.2	2	12	$\bar{2}$	27.5	-24.8
0	5	1	19.3	17.3	4	12	$\bar{4}$	20.5	30.7
0	7	1	18.5	16.3	1	14	$\bar{1}$	42.0	-40.5
0	9	1	65.2	67.2	2	14	$\bar{2}$	32.0	-24.3
0	11	1	36.0	41.4	3	14	$\bar{3}$	< 3.0	5.6
0	13	1	13.9	12.7	1	16	$\bar{1}$	43.0	-44.9
0	15	1	25.0	26.1					

splitting of fundamental frequencies associated with the  $\text{SO}_4$  ion. The lower symmetry was not revealed by Atoji & Rundle from their neutron diffraction results but the re-refining of their observations shows (Table 5) that it was present in their data. The reduced symmetry probably arises because of the different environment surrounding the two crystallographically distinct sulphate oxygens O(1) and O(11). The symmetry-related pair of atoms O(1) and O(1') are each hydrogen bonded to two water molecules towards which they face whereas the symmetry-related pair O(11) and O(11') form no hydrogen bonds but have close  $\text{Ca}^{2+}$  neighbours. Atoms O(1), O(11) and O(1'), O(11') are drawn closer together and atoms O(1), O(1') and O(11), O(11') spread slightly apart with respect to the values found by Atoji & Rundle to produce the observed distortion from tetrahedral symmetry.

Each calcium atom in gypsum is surrounded by three pairs of oxygen atoms [one of O(1), O(1') and two of O(11), O(11')] belonging to  $\text{SO}_4$  groups and by a pair of water molecules (Fig. 1) with oxygen atoms O( $W$ ), O( $W'$ ). For atoms O(1) and O(11) of the same  $\text{SO}_4$  group the Ca-O(1) pair and the Ca-O(11) pair have bond lengths that are significantly shorter (Table 5) than those found by Atoji & Rundle. This arises from the shortening of the O(1)-O(11) bond distance and the decreasing of the O(1)-S-O(11) bond angle. The second shorter Ca-O(11) bond length (Table 5) links the Ca of the one cell to the symmetry-related pair of O(11) and O(11') atoms in the next cell and this is slightly lengthened, as is the distance between Ca and the O( $W$ ) atoms in the same cell with respect to the values found by Atoji & Rundle.

The changes in the  $x$  and  $z$  parameters of atoms O(1) and O( $W$ ) (Table 1) from those found by Atoji & Rundle affect the hydrogen-bond distances given by O( $W$ )-O(1) [or O(1')] through H(1) and O( $W'$ )-O(1) [or O(1')] through H(2) (Fig. 1). These distances, as well as the two distances O( $W$ )-H(1) and O( $W$ )-H(2) in each water molecule, are non-equivalent and in principle need not be of the same length (Seidl, Knop & Falk, 1969). In fact Seidl, Knop & Falk have shown that appreciable asymmetry of the water molecule is indicated by the large splittings in gypsum of the three isolated HDO frequencies:  $90\text{ cm}^{-1}$  for O-H stretching,  $65\text{ cm}^{-1}$  for O-D stretching, and  $14\text{ cm}^{-1}$  for the bending. They argue that it would be necessary to determine the positional parameters of O( $W$ ) and O(1) with standard deviations of the order of  $0.001\text{ \AA}$  to verify the asymmetry found by the spectroscopic results. The new X-ray results indicate that the two distances O( $W$ )-O(1) [or O(1')] through H(1) and O( $W$ )-O(1) [or O(1')] through H(2) of the water molecule differ by  $0.080 \pm 0.003\text{ \AA}$  (Table 5) while the re-refined neutron data indicate a difference of  $0.065 \pm 0.018\text{ \AA}$ , both results being statistically significant. The original refinement of Atoji & Rundle gave no difference in these bond lengths.

If the true hydrogen positions are taken as those

Table 5. *Interatomic distances and bond angles*

	X-ray, isotropic and anisotropic This study	Neutron, isotropic	
		Atoji & Rundle	Re-refined
<b>Sulphate ion</b>			
S—O(1)	1.457 ± 0.001 Å	1.497 ± 0.014 Å	1.483 ± 0.014 Å
S—O(11)	1.461 ± 0.002	1.479 ± 0.010	1.468 ± 0.014
O(1)—O(1')	2.409 ± 0.002	2.442 ± 0.013	2.452 ± 0.020
O(1)—O(11) = O(1')—O(11')	2.323 ± 0.002	2.421 ± 0.014	2.370 ± 0.013
O(11)—O(11')	2.407 ± 0.003	2.415 ± 0.010	2.398 ± 0.018
O(1)—O(11') = O(1')—O(11)	2.414 ± 0.002	2.438 ± 0.011	2.434 ± 0.011
O(1)—S—O(1')	111.5 ± 0.1°	109.4 ± 1.4°	111.5 ± 1.3°
O(11)—S—O(11')	111.0 ± 0.1	109.4 ± 1.0	109.5 ± 1.3
O(1)—S—O(11) = O(1')—S—O(11')	105.5 ± 0.1	108.9 ± 1.3	106.8 ± 0.5
O(1)—S—O(11') = O(1')—S—O(11)	111.7 ± 0.1	110.1 ± 1.2	111.1 ± 0.6
<b>Water molecule</b>			
O(W)—H(1)	0.962 ± 0.012 Å	1.002 ± 0.030 Å	0.981 ± 0.040 Å
O(W)—H(2)	0.944 ± 0.021	0.981 ± 0.024	0.939 ± 0.023
H(1)—H(2)	1.531 ± 0.030	1.580 ± 0.036	1.531 ± 0.030
H(1)—O(W)—H(2)	106.9 ± 1.0°	105.6 ± 4.5°	105.7 ± 1.8°
<b>Hydrogen bond</b>			
O(W)—O(1) [or O(1') through H(1)]	2.816 ± 0.002 Å	2.816 ± 0.014 Å	2.799 ± 0.013 Å
O(W)—O(1) [or O(1') through H(2)]	2.896 ± 0.002	2.824 ± 0.015	2.865 ± 0.011
O(1)—H(1) = O(1')—H(1)	1.860 ± 0.013	1.814	1.826 ± 0.038
O(1)—H(2) = O(1')—H(2)	1.955 ± 0.006	1.845	1.932 ± 0.019
O(1)—H(1)—O(W) = O(1')—H(1)—O(W)	172.3 ± 1.0°	177.7°	171.0 ± 1.8°
O(1)—H(2)—O(W) = O(1')—H(2)—O(W)	174.1 ± 1.0	178.1	172.3 ± 2.1
<b>Calcium—oxygen distances</b>			
Ca—O(1)	2.528 ± 0.002 Å	2.566 Å	2.552 ± 0.009 Å
Ca—O(11)	2.544 ± 0.002	2.588	2.548 ± 0.010
Ca—O(11)	2.378 ± 0.001	2.353	2.357 ± 0.012
Ca—O(W)	2.380 ± 0.001	2.341	2.383 ± 0.014

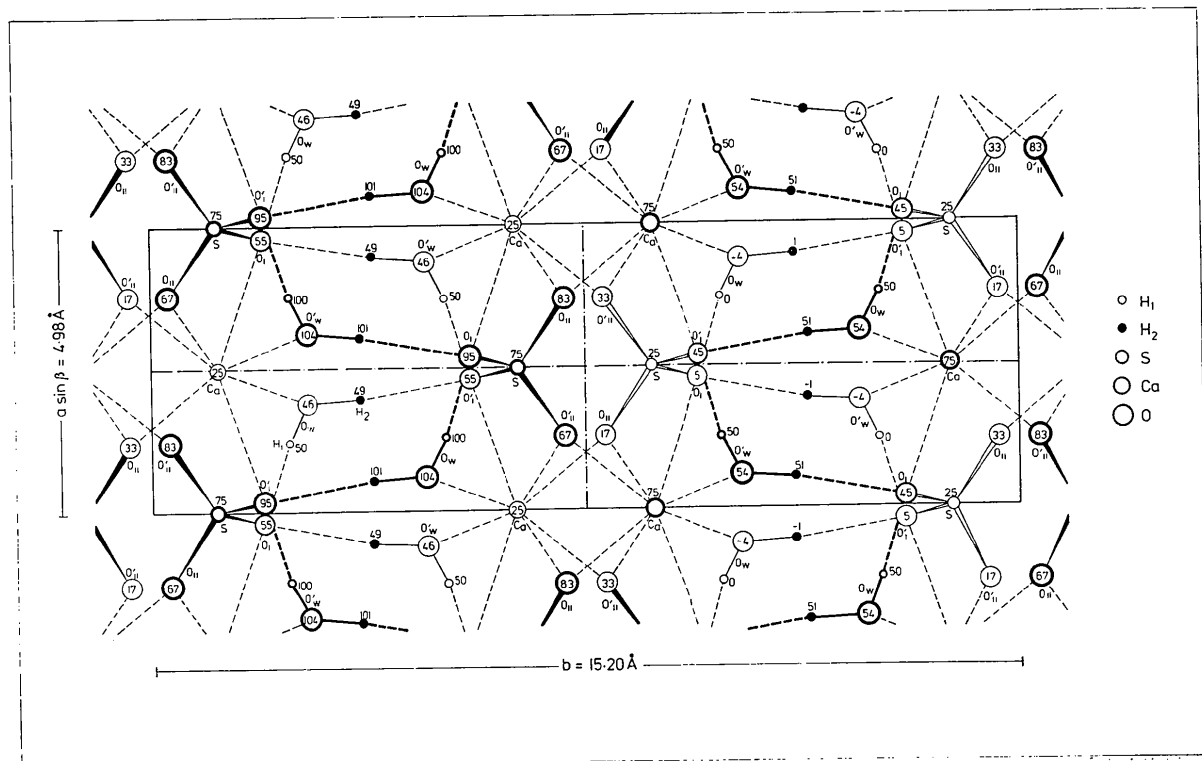


Fig. 1. One-half of the unit cell of gypsum projected on the (001) plane [after Atoji &amp; Rundle (1958)]. The fractional coordinates have been multiplied by 100.

given from the re-refined neutron data (Table 1) and if they are combined with the  $O(W)$  parameters given by the X-ray refinement, then the  $O(W)$ -H(1) distance is 0.962 Å, and the  $O(W)$ -H(2) distance 0.944 Å, the H(1)-O(W)-H(2) bond angle 106.9°, and the H(1)-H(2) distance 1.531 Å (Table 5). The  $O(W)$ -H and  $O(W)$ -H-O(1) distances correlate well with those from hydrates given in Fig. 8 of Atoji & Rundle. The location of the two hydrogen atoms is essentially along the lines joining  $O(W)$ -O(1) and  $O(W)$ -O(1') (Table 5 and Fig. 1) as reported earlier by Atoji & Rundle. Thus the shorter  $O(W)$ -H(2) distance is associated with the longer  $O(W)$ -H(2)-O(1) [or O(1')] distance and the longer  $O(W)$ -H(1) with the shorter  $O(W)$ -H(1)-O(1) [or O(1')] as expected. The ratio of the distance  $O(W)$ -H(2) to the distance  $O(W)$ -H(1) is 0.979 which can be compared with the ratio of 0.966 reported by Seidl, Knop & Falk (1969) as being derivable by R. W. Green from spectroscopic data.

In the attempts to re-refine the neutron-diffraction data with mixed isotropic-anisotropic temperature factors the hydrogen positional parameters, after one cycle of refinement, gave higher values for the  $O(W)$ -H(1) (1.042 Å) and  $O(W)$ -H(2) (0.937 Å) distances, using the  $O(W)$  positional parameters from the neutron refinement as also when using the X-ray positional parameters for  $O(W)$  [ $O(W)$ -H(1) 1.022 Å and  $O(W)$ -H(2) 0.951 Å]. Thus, although the hydrogen positional parameters are known with less certainty than the others, those obtained in the present work, when combined with  $O(W)$  positions from either neutron or X-ray data, yield values of bond lengths consistent with two hydrogen bonds of different length and strength.

In the structure of gypsum the water molecules form two parallel sheets in the direction of the prominent cleavage of the crystal (Fig. 1), and the H(2) atoms are all located within the double layer whereas the H(1) atoms are always outside it. This means that an electrical repulsion between the sheets would lengthen hydrogen bonds involving the H(2) atom relative to those involving the H(1) atom, as is observed.

In the neutron-diffraction refinement of Atoji & Rundle the parameters were refined with isotropic temperature factors, since it was observed from Fourier methods that the thermal motions were isotropic or nearly isotropic. The root-mean square amplitudes of the thermal motion along the principal axes set out

in Table 6 show that the thermal motions of Ca and S are nearly isotropic, but those of O(1), O(11) and  $O(W)$  show some anisotropy, particularly  $O(W)$ . The locations of the thermal ellipsoids for Ca and S are very similar and are shown in Fig. 2 with respect to the atom projections on the  $ac$  plane. Greatest thermal movement of

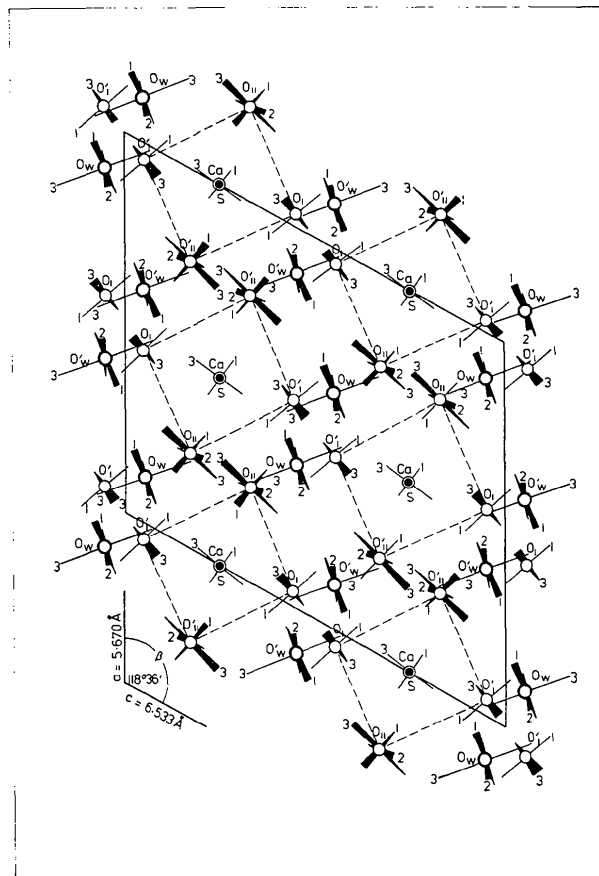


Fig. 2. Position of principal axes of thermal ellipsoids (1, 2 and 3) with respect to the  $ac$  plane of gypsum. The lengths of the axes are proportional to the magnitude of the r.m.s. amplitudes of thermal motion (Table 6) projected into the  $ac$  plane. The thicknesses of the wedges indicate the magnitude of the angle of inclination ( $I$ ) of the principal axes to the  $ac$  plane. The directions that the wedges point to are the angles  $\phi$  (Table 6) that the principal axes make with the  $a$  axis. For atoms O(1) and O(1') the principal axis 2 nearly superimposes on axis 3 and is not shown. Its inclination to the plane is opposite to that of axis 3.

Table 6. Root-mean-square (r.m.s.) amplitudes of thermal motion along the principal axes, angle of principal axes to the  $a$  axis ( $\phi$ ) and inclination of principal axes to  $ac$  plane ( $I$ )

(Numbers in parentheses are standard deviations in the least significant digit.)

	Axis 1			Axis 2			Axis 3		
	$\phi$	$I$	r.m.s.	$\phi$	$I$	r.m.s.	$\phi$	$I$	r.m.s.
Ca	45°	0°	0.092 (1) Å	90°	0°	0.105 (1) Å	135°	0°	0.108 (1) Å
S	37	0	0.077 (2)	—	90	0.091 (1)	127	0	0.100 (1)
O(1)	50	0	0.118 (2)	320°	32	0.118 (2)	142	58	0.140 (2)
O(11)	42	9	0.093 (2)	293	66	0.118 (2)	135	22	0.154 (2)
$O(W)$	157	26	0.121 (3)	347	64	0.150 (2)	249	3	0.191 (2)



Ca and S takes place approximately in the direction of the  $c$  axis. The intermediate axis of these ellipsoids lies in the direction of the  $b$  axis. The directions of the ellipsoids for O(1) and O(11) are similar to those for Ca and S except that the intermediate axis is inclined to the  $b$  axis (Table 6). For atom O( $W$ ) the directions of the ellipsoid are different from those of the other atoms and the greatest thermal motion is more nearly in line with the least thermal motion of atoms O(1), O(11), Ca and S. Its intermediate axis is inclined to the  $b$  axis (Table 6).

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## Electron Diffraction and the Structure of $\alpha$ -N<sub>2</sub>

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Electron diffraction patterns have been obtained from annealed single crystals of  $\alpha$ -N<sub>2</sub>. The {100} patterns were found to be particularly suitable for resolving the controversy as to whether  $\alpha$ -N<sub>2</sub> has the  $Pa3$  or  $P2_13$  structure. A quantitative comparison of these patterns with calculated intensities was used to set a (conservative) upper limit of 0.05 Å for the displacement of the centre of the N<sub>2</sub> molecule from the  $Pa3$  structure. This and other more qualitative evidence strongly suggests that the structure is in fact  $Pa3$ . Twins on {111} planes were observed and were shown to give rise to some reflexions which are forbidden by  $Pa3$  symmetry. The possibility is discussed that twinning could explain a previous X-ray observation of isolated reflexions which were not consistent with the  $Pa3$  structure.

### 1. Introduction

The exact structure of the low temperature  $\alpha$ -phase of solid nitrogen has been under discussion since the first X-ray work of Vegard (1929) and Ruhemann (1932). The structure proposed by Ruhemann was  $Pa3$  cubic, in which the centres of the N<sub>2</sub> molecules are on a f.c.c.

lattice, but each molecule points in a different  $\langle 111 \rangle$  direction. The Vegard structure is essentially similar but the centre of a molecule is displaced by a small distance,  $r$ , along the  $\langle 111 \rangle$  direction parallel to the axis of that molecule. This displacement lowers the symmetry to  $P2_13$ , which is a non-centrosymmetric structure.

Many diffraction experiments have been performed in an attempt to decide which structure is correct. X-ray powder patterns (Bolz, Boyd, Mauer & Peiser,

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